

## ON THE ALGEBRA OF 3-DIMENSIONAL *ES*-MANIFOLD

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ABSTRACT. The manifold  $*g-ESX_n$  is a generalized  $n$ -dimensional Riemannian manifold on which the differential geometric structure is imposed by the unified field tensor  $*g^{\lambda\nu}$  through the *ES*-connection which is both Einstein and semi-symmetric. The purpose of the present paper is to study the algebraic geometric structures of 3-dimensional  $*g-ESX_3$ . Particularly, in 3-dimensional  $*g-ESX_3$ , we derive a new set of powerful recurrence relations in the first class.

### 1. Introduction

In Appendix *II* to his last book Einstein([4], 1950) proposed a new unified field theory that would include both gravitation and electromagnetism. Although the intent of this theory is physical, its exposition is mainly geometrical. It may be characterized as a set of geometrical postulates for the space time  $X_4$ . Characterizing Einstein's unified field theory as a set of geometrical postulates for  $X_4$ , Hlavatý([5], 1957) gave its mathematical foundation for the first time. Since then Hlavatý and number of mathematicians contributed for the development of this theory and obtained many geometrical consequences of these postulates.

The main purpose of the present paper is to study the algebraic geometric properties of 3-dimensional  $*g-ESX_3$  in the first class. In

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particular, we derive a powerful recurrence relations in 3-dimensional  $*g - ESX_3$ .

## 2. Preliminaries

This section is a brief collection of basic concepts, results, and notations needed in subsequent considerations. They are due to Hwang ([1], 1995), Chung ([2], 1963), and Mishra([6], 1959) mostly due to [3].

### (a) $n$ -dimensional $*g$ -unified field theory

Corresponding to the Einstein's  $n - g$ -UFT, our  $n - *g$ -UFT is based on the following three principles.

*Principle A.* Let  $X_n$  be an  $n$ -dimensional generalized Riemannian manifold referred to a real coordinate system  $x^\nu$ , which obeys the coordinate transformations  $x^\nu \rightarrow x'^\nu$  for which

$$(2.1) \quad \det\left(\frac{\partial x'}{\partial x}\right) \neq 0$$

In  $n - g - UFT$  the manifold  $X_n$  is endowed with a real nonsymmetric tensor  $g_{\lambda\mu}$ , which may be decomposed into its symmetric part  $h_{\lambda\mu}$  and skew-symmetric part  $k_{\lambda\mu}$ :

$$(2.2) \quad g_{\lambda\mu} = h_{\lambda\mu} + k_{\lambda\mu}$$

where

$$(2.3) \quad \mathfrak{g} = \det(g_{\lambda\mu}) \neq 0, \quad \mathfrak{h} = \det(h_{\lambda\mu}) \neq 0, \quad \mathfrak{k} = \det(k_{\lambda\mu})$$

In  $n - *g - UFT$  the algebraic structure on  $X_n$  is imposed by the basic real tensor  $*g^{\lambda\nu}$  defined by

$$(2.4) \quad g_{\lambda\mu} *g^{\lambda\nu} = g_{\mu\lambda} *g^{\nu\lambda} = \delta_\mu^\nu$$

It may be also decomposed into its symmetric part  $*h^{\lambda\nu}$  and skew-symmetric part  $*k^{\lambda\nu}$ :

$$(2.5) \quad *g^{\lambda\nu} = *h^{\lambda\nu} + *k^{\lambda\nu}$$

Since  $\det(*h^{\lambda\nu}) \neq 0$ , we may define a unique tensor  $*h_{\lambda\mu}$  by

$$(2.6) \quad *h_{\lambda\mu} *h^{\lambda\nu} = \delta_\mu^\nu$$

In  $n - *g$ -UFT we use both  $*h^{\lambda\nu}$  and  $*h_{\lambda\mu}$  as tensors for raising and/or lowering indices of all tensors in  $X_n$  in the usual manner. We then have

$$(2.7) \quad *k_{\lambda\mu} = *k^{\rho\sigma} *h_{\lambda\rho} *h_{\mu\sigma}, \quad *g_{\lambda\mu} = *g^{\rho\sigma} *h_{\lambda\rho} *h_{\mu\sigma}$$

so that

$$(2.8) \quad *g_{\lambda\mu} = *h_{\lambda\mu} + *k_{\lambda\mu}$$

*Principle B* The differential geometric structure on  $X_n$  is imposed by the tensor  $*g^{\lambda\nu}$  by means of a connection  $\Gamma_{\lambda}^{\nu\mu}$  defined by a system of equations

$$(2.9) \quad D_{\omega} *g^{\lambda\nu} = -2S_{\omega\alpha}^{\nu} *g^{\lambda\alpha}$$

Here  $D_{\omega}$  denotes the symbol of the covariant derivative with respect to  $\Gamma_{\lambda}^{\nu\mu}$  and  $S_{\lambda\mu}^{\nu}$  is the torsion tensor of  $\Gamma_{\lambda}^{\nu\mu}$ . Under certain conditions the system (2.9) admits a unique solutions  $\Gamma_{\lambda}^{\nu\mu}$ .

*Principle C* In order to obtain  $*g^{\lambda\nu}$  involved in the solution for  $\Gamma_{\lambda}^{\nu\mu}$  certain conditions are imposed. These conditions may be condensed to

$$(2.10) \quad S_{\lambda} = S_{\lambda\alpha}^{\alpha} = 0, \quad R_{[\mu\lambda]} = \partial_{[\mu} Y_{\lambda]}, \quad R_{(\mu\lambda)} = 0$$

where  $Y_{\lambda}$  is an arbitrary vector, and  $R_{\omega\mu\lambda}^{\nu}$  are the curvature tensors of  $X_n$  defined by

$$(2.11) \quad R_{\omega\mu\lambda}^{\nu} = 2(\partial_{[\mu} \Gamma_{|\lambda|}^{\nu\omega]} + \Gamma_{\alpha}^{\nu}{}_{[\mu} \Gamma_{|\lambda|}^{\alpha\omega]}), \quad R_{\mu\lambda} = R_{\alpha\mu\lambda}^{\alpha}$$

### (b) Some notations and results

The following quantities are frequently used in our further considerations:

$$(2.12) \quad *g = \det(*g_{\lambda\mu}), \quad *h = \det(*h_{\lambda\mu}), \quad *k = \det(*k_{\lambda\mu})$$

$$(2.13) \quad *g = \frac{*g}{*h}, \quad *k = \frac{*k}{*h}.$$

$$(2.14) \quad K_p = *k_{[\alpha_1}^{\alpha_1} *k_{\alpha_2}^{\alpha_2} \dots *k_{\alpha_p]}^{\alpha_p}, \quad (p = 0, 1, 2, \dots).$$

$$(2.15) \quad {}^{(0)}*k_{\lambda}^{\nu} = \delta_{\lambda}^{\nu}, \quad {}^{(p)}*k_{\lambda}^{\nu} = *k_{\lambda}^{\alpha} {}^{(p-1)}*k_{\alpha}^{\nu} \quad (p = 1, 2, \dots).$$

In  $X_n$  it was proved in [2] that

$$(2.16) \quad K_0 = 1, \quad K_n = *k \text{ if } n \text{ is even, and } K_n = 0 \text{ if } n \text{ is odd.}$$

$$(2.17) \quad {}^*g = 1 + K_2 + \cdots + K_{n-\sigma}.$$

$$(2.18) \quad \sum_{s=0}^{n-\sigma} K_s \, {}^{(n-s)*}k_{\lambda}{}^{\nu} = 0 \quad (p = 0, 1, 2, \dots).$$

We also use the following useful abbreviations for an arbitrary tensor  $T_{\dots}$  for  $p = 1, 2, 3, \dots$ :

$$(2.19) \quad {}^{(p)}T_{\dots} = {}^{(p-1)}k^{\nu}{}_{\alpha} T_{\dots}^{\alpha \dots}.$$

### (c) $n$ -dimensional $ES$ manifold $n - {}^*g$ -UFT

In this subsection, we display an useful representation of the  $ES$  connection in  $n - {}^*g$ -UFT.

DEFINITION 2.1. A connection  $\Gamma_{\lambda}{}^{\nu}{}_{\mu}$  is said to be *semi-symmetric* if its torsion tensor  $S_{\lambda\mu}{}^{\nu}$  is of the form

$$(2.20) \quad S_{\lambda\mu}{}^{\nu} = 2\delta_{[\lambda}^{\nu} X_{\mu]}.$$

for an arbitrary non-null vector  $X_{\mu}$ .

A connection which is both semi-symmetric and Einstein is called an  $ES$  connection. An  $n$ -dimensional generalized Riemannian manifold  $X_n$ , on which the differential geometric structure is imposed by  ${}^*g^{\lambda\nu}$  by means of an  $ES$  connection, is called an  $n$ -dimensional  ${}^*g$ - $ES$  manifold. We denote this manifold by  ${}^*g$ - $ESX_n$  in our further considerations.

THEOREM 2.2. Under the condition (2.20), the system of equations (2.9) is equivalent to

$$(2.21) \quad \Gamma_{\lambda}{}^{\nu}{}_{\mu} = {}^* \left\{ \begin{array}{c} \nu \\ \lambda\mu \end{array} \right\} + U^{\nu}{}_{\lambda\mu} + 2\delta_{[\lambda}^{\nu} X_{\mu]}.$$

where

$$(2.22) \quad U^{\nu}{}_{\lambda\mu} = -{}^*h_{\lambda\mu}{}^{(2)} X^{\nu}$$

*Proof.* Substituting (2.20) for  $S_{\lambda\mu}{}^{\nu}$  into (2.9), we have the representation (2.21).  $\square$

### 3. Recurrence relations in $*g - ESX_3$

In this section we derive several powerful recurrence relations, establishing a nonholonomic frame in  $*g - ESX_3$ .

DEFINITION 3.1. The tensors  $*g_{\lambda\mu}$  is said to be

- (1) of the first class, if  $K_{n-\sigma} \neq 0$
- (2) of the second class with  $j$ th category ( $j \geq 1$ ), if

$$(3.1) \quad K_{2j} \neq 0, \quad K_{2j+2} = K_{2j+4} = \dots = K_{n-\sigma} = 0$$

- (3) of the third class, if  $K_2 = K_4 = \dots = K_{n-\sigma} = 0$

In  $3 - *g$ -UFT, we have only two classes; namely the first class when  $K_2 \neq 0$  and the third class when  $K_2 = 0$ . In  $3 - *g$ -UFT, the relation (2.17) gives

$$(3.2) \quad *g = 1 + K_2$$

#### A. Basic vectors in the first class

The eigenvalues  $M$  and the corresponding eigenvectors  $A^\nu$  in  $*g - ESX_n$ , defined by

$$(3.3) \quad MA^\nu = *k_\mu{}^\nu A^\mu, \quad (M : \text{a scalar}).$$

are called *basic scalars* and *basic vectors*, respectively.

THEOREM 3.2. The basic scalars in  $*g - ESX_3$  may be given by

$$(3.4) \quad M_1 = -M_2 = \sqrt{-K_2}, \quad M_3 = 0$$

*Proof.* In  $3 - *g$ -UFT, the characteristic equation is reduced to

$$(3.5) \quad M(M^2 + K_2) = 0$$

from which our assertion follows. □

THEOREM 3.3. There are three linearly independent basic vectors  $A_1^\nu, A_2^\nu, A_3^\nu$  and they have the following properties:

- (a) They are defined up to an arbitrary factor of proportionality.
- (b)  $A_1^\nu, A_2^\nu$  are null vectors, while  $A_3^\nu$  is non-null.
- (c)  $A_3^\nu$  is perpendicular to  $A_1^\nu$  and  $A_2^\nu$ .
- (d)  $A_1^\nu$  and  $A_2^\nu$  satisfy the condition  $h_{\lambda\mu} A_1^\lambda A_2^\mu \neq 0$ .

*Proof.* Since the basic scalars  $M_x^i$  are all distinct, (3.3) admits three linearly independent basic vectors  $A_x^\nu$  which are defined up to an arbitrary factor of proportionality. The first half of statement (b) is a consequence of (3.3), (3.4), and

$$(3.6) \quad M_x^* h_{\lambda\mu} A_x^\lambda A_x^\mu = {}^*k_{\lambda\mu} A_x^\lambda A_x^\mu = 0, \quad (x = 1, 2)$$

Since  $M_3 + M_x \neq 0$ ,  $(x = 1, 2)$ , statement (c) follows from (3.4) as in the following way:

$$(3.7) \quad M_x^* h_{\lambda\mu} A_x^\lambda A_x^\mu = {}^*k_{\lambda\mu} A_x^\lambda A_x^\mu = -M_3^* h_{\mu\lambda} A_x^\lambda A_x^\mu$$

In order to prove statement (d), consider a conic  $C$  with equation  ${}^*h_{\lambda\mu} A^\lambda A^\mu = 0$  on a projective plane  $P_2$ . In virtue of statement (b),  $A_1^\nu$  and  $A_2^\nu$  are two different points on  $C$  while  ${}^*h_{\lambda\mu} A_1^\lambda = A_{1\mu}$  is the tangent line to  $C$  at  $A_1^\nu$ . Since  $\det({}^*h_{\lambda\mu}) \neq 0$ ,  $C$  is non-degenerate. Consequently  ${}^*h_{\lambda\mu} A_1^\lambda = A_{1\mu}$  and  $A_2^\mu$  are not incident; that is,  ${}^*h_{\lambda\mu} A_1^\lambda A_2^\mu \neq 0$ .  $\square$

REMARK 3.4. We may choose the factor of proportionality mentioned in Theorem (3.3) as

$$(3.8) \quad {}^*h_{12} = {}^*h_{33} = 1$$

## B. Nonholonomic frame of reference in the first class

In the first class, we have a set of 3 linearly independent basic vectors  $A_x^\nu$ ,  $(i = 1, 2, 3)$  and a unique reciprocal set  $A_\lambda^i$ ,  $(i = 1, 2, 3)$  satisfying

$$(3.9) \quad A_\lambda^j A_\lambda^i = \delta_i^j, \quad A_\lambda^i A_\lambda^\nu = \delta_\lambda^\nu$$

With these two set of vectors, we may construct a nonholonomic frame of reference as follows;

DEFINITION 3.5. If  $T_{\lambda\dots}^{\nu\dots}$  are holonomic components of a tensor, then its nonholonomic components are defined by

$$(3.10) \quad T_{j\dots}^{i\dots} = T_{\lambda\dots}^{\nu\dots} A_\nu^i \cdots A_j^\lambda \cdots$$

An easy inspection shows that

$$(3.11) \quad T_{\lambda \dots}^{\nu \dots} = T_{j \dots}^{i \dots} A_i^{\nu} \cdots A_{\lambda}^j \cdots$$

**THEOREM 3.6.** *The nonholonomic components  ${}^*h_{ij}$  and  ${}^*h^{ij}$  may be given by the matrix equation*

$$(3.12) \quad ({}^*h_{ij}) = ({}^*h^{ij}) = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

*Proof.* (3.12) is a direct result of (3.8) and theorem (3.3) □

**THEOREM 3.7.** *We have*

$$(3.13) \quad A_i^{\nu} = A_{\lambda}^j {}^*h_{ij} {}^*h^{\lambda\nu}, \quad A_{\lambda}^j = A_i^{\nu} {}^*h^{ij} {}^*h_{\lambda\nu}$$

so that

$$(3.14) \quad \begin{matrix} A_1^{\nu} = A^2_{\nu} & A_2^{\nu} = A^1_{\nu} & A_3^{\nu} = A^3_{\nu} \\ A_{\lambda}^1 = A_{2\lambda} & A_{\lambda}^2 = A_{1\lambda} & A_{\lambda}^3 = A_{3\lambda} \end{matrix}$$

*Proof.* In virtue of (2.6), (3.9) and (3.10), the first relation of (3.13) follows as in the following way;

$$(3.15) \quad A_{\lambda}^j {}^*h_{ij} {}^*h^{\lambda\nu} = A_{\lambda}^j ({}^*h_{\alpha\beta} A_i^{\alpha} A_j^{\beta}) {}^*h^{\lambda\nu} = A_i^{\alpha} ({}^*h_{\alpha\beta} {}^*h^{\lambda\nu}) \delta_{\lambda}^{\beta} = A_i^{\nu}$$

(3.14) follows from (3.13) and (3.12). □

**THEOREM 3.8.** *The nonholonomic components of  ${}^{(p)*}k_{\lambda}{}^{\nu}$ ,  ${}^{(p)*}k_{\lambda\mu}$  and  ${}^{(p)*}k^{\lambda\nu}$  are given by*

$$(3.16) \quad {}^{(p)*}k_x{}^i = M_x^p \delta_x^i$$

$$(3.17) \quad {}^{(p)*}k_{xi} = M_x^{p*} h_{xi}$$

$$(3.18) \quad {}^{(p)*}k^{xi} = M_x^{p*} h^{xi}$$

*Proof.* Let  $A^{\nu}_x$  be the basic vector corresponding to the basic scalar  $M$ . Then from (3.3), we have

$$(3.19) \quad {}^{(p)*}k_{\lambda}^{\nu} A^{\lambda}_x = M^p A^{\nu}_x \quad (p = 0, 1, 2, \dots)$$

(3.16) follows immediately by multiplying  $A^i_{\nu}$  to both sides of (3.19). The remaining relations may be obtained from (3.16) by lowering and/or raising indices.  $\square$

In the following theorem, we express the components of tensors  $*h_{\lambda\mu}$ ,  ${}^{(p)*}k_{\lambda}^{\nu}$ ,  ${}^{(p)*}k_{\lambda\mu}$ ,  ${}^{(p)*}k^{\lambda\nu}$  in terms of basic vectors:

**THEOREM 3.9.** *The representation of  $*h_{\lambda\mu}$ ,  ${}^{(p)*}k_{\lambda}^{\nu}$ ,  ${}^{(p)*}k_{\lambda\mu}$ ,  ${}^{(p)*}k^{\lambda\nu}$  in terms of basic vectors may be given by*

$$(3.20) \quad *h_{\lambda\mu} = 2A^1_{(\lambda} A^2_{\mu)} + A^3_{\lambda} A^3_{\mu}$$

$$(3.21) \quad {}^{(p)*}k_{\lambda}^{\nu} = M^p (A^1_{\lambda} A^{\nu}_1 + (-1)^p A^2_{\lambda} A^{\nu}_2)$$

$$(3.22) \quad {}^{(p)*}k_{\lambda\mu} = \begin{cases} 2M^p A^1_{(\lambda} A^2_{\mu)}, & \text{if } p \text{ is even} \\ 2M^p A^1_{[\lambda} A^2_{\mu]}, & \text{if } p \text{ is odd} \end{cases}$$

$$(3.23) \quad {}^{(p)*}k^{\lambda\nu} = \begin{cases} 2M^p A^{(\lambda}_1 A^{\nu)}_2, & \text{if } p \text{ is even} \\ 2M^p A^{[\lambda}_1 A^{\nu]}_2, & \text{if } p \text{ is odd} \end{cases}$$

*Proof.* The representations (3.20) – (3.23) follow from (3.11) in virtue of (3.4), (3.12) and (3.14).  $\square$

### C. Recurrence relations in the first class

In this subsection we derive several recurrence relations.

**THEOREM 3.10.** *In the first class, the tensor  $T_{\omega\mu\nu}$ , skew-symmetric in the first two indices, satisfies*

$$(3.24) \quad {}^{(pq)r}T_{\omega\mu\nu} = \sum_{x,y,z} T_{xyz} M^p_x M^q_y M^r_z A^x_{\omega} A^y_{\mu} A^z_{\nu}$$

$$(3.25) \quad T_{\nu[\omega\mu]}^{r(pq)} = \sum_{x,y,z} T_{x[yz]} M_x^{(p)} M_y^{(q)} M_z^r \overset{x}{A}_\nu \overset{y}{A}_\omega \overset{z}{A}_\mu$$

*Proof.* In virtue of (3.11) and (3.16), our assertion (3.24) may be derived as

$$(3.26) \quad \begin{aligned} T_{\omega\mu\nu}^{(pq)r} &= \sum_{x,y,z} T_{xyz}^{(pq)r} \overset{x}{A}_\omega \overset{y}{A}_\mu \overset{z}{A}_\nu \\ &= \sum_{x,y,z} \frac{1}{2} \left( (p)^* k_x^i (q)^* k_y^j + (q)^* k_x^i (p)^* k_y^j \right) (r)^* k_z^k \overset{x}{A}_\omega \overset{y}{A}_\mu \overset{z}{A}_\nu \\ &= \frac{1}{2} \sum_{x,y,z} T_{xyz} (M_x^p M_y^q + M_x^q M_y^p) M_z^r \overset{x}{A}_\omega \overset{y}{A}_\mu \overset{z}{A}_\nu \end{aligned}$$

The second relation may be proved similarly. □

**THEOREM 3.11.** *The main recurrence relation in the first class is*

$$(3.27) \quad (p+3)^* k_\lambda^\nu = -K_2^{(p+1)*} k_\lambda^\nu, \quad (p = 0, 1, 2, \dots)$$

*Proof.* Let  $M_x$  be a basic scalar. In  $3 - *g - ESX_3$ , the characteristic equation is

$$(3.28) \quad \sum_{f=0}^2 K_f M_x^{3-f} = 0$$

Multiplying  $\delta_x^i$  to both sides of (3.28) and making use of (3.16), we have

$$(3.29) \quad \sum_{f=0}^2 K_f (3-f)^* k_x^i = 0$$

whose holonomic form is

$$(3.30) \quad \sum_{f=0}^2 K_f (3-f)^* k_\lambda^\alpha = 0$$

The relation (3.27) immediately follows by multiplying  $(p)^* k_\alpha^\nu$  to both sides of (3.30). □

**THEOREM 3.12.** *The basic scalars  $M_x$  satisfy*

$$(3.31) \quad M_x M_y (M_x + M_y), \quad (x \neq y)$$

$$(3.32) \quad M_x M_y (M_x M_y - K_2), \quad (x \neq y)$$

*Proof.* These relations follow easily from (3.4). □

**THEOREM 3.13.** (*Recurrence relations in the first class*) If  $T_{\omega\mu\nu}$  is a tensor skew-symmetric in the first two indices, then the following recurrence relations hold in the first class of  $3 - *g - ESX_3$  :

$$(3.33) \quad \begin{matrix} (12)r \\ T \end{matrix} = 0, \quad \begin{matrix} 22r \\ T \end{matrix} = K_2 \begin{matrix} 11r \\ T \end{matrix}$$

$$(3.34) \quad \begin{matrix} r(12) \\ T \end{matrix} \nu[\omega\mu] = 0, \quad \begin{matrix} r22 \\ T \end{matrix} \nu[\omega\mu] = K_2 \begin{matrix} r11 \\ T \end{matrix} \nu[\omega\mu]$$

*Proof.* The above relations are consequences of (3.24), (3.25) and (3.31), (3.32). For example, the second relation of (3.33) may be proved as in the following way:

$$(3.35) \quad \begin{aligned} \begin{matrix} 22r \\ T \end{matrix} \omega\mu\nu &= \sum_{x,y,z} T_{xyz} M_x^2 M_y^2 M_z^r A_\omega^x A_\mu^y A_\nu^z \\ &= K_2 \sum_{x,y,z} T_{xyz} M_x M_y M_z^r A_\omega^x A_\mu^y A_\nu^z \\ &= K_2 \begin{matrix} 11r \\ T \end{matrix} \omega\mu\nu = K_2 \begin{matrix} 11r \\ T \end{matrix} \end{aligned}$$

□

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